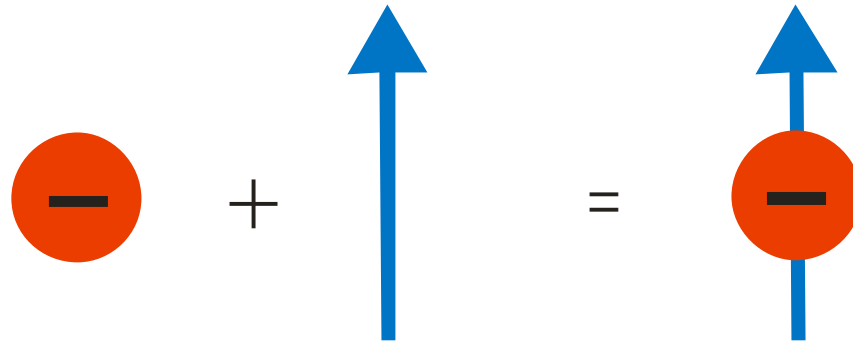


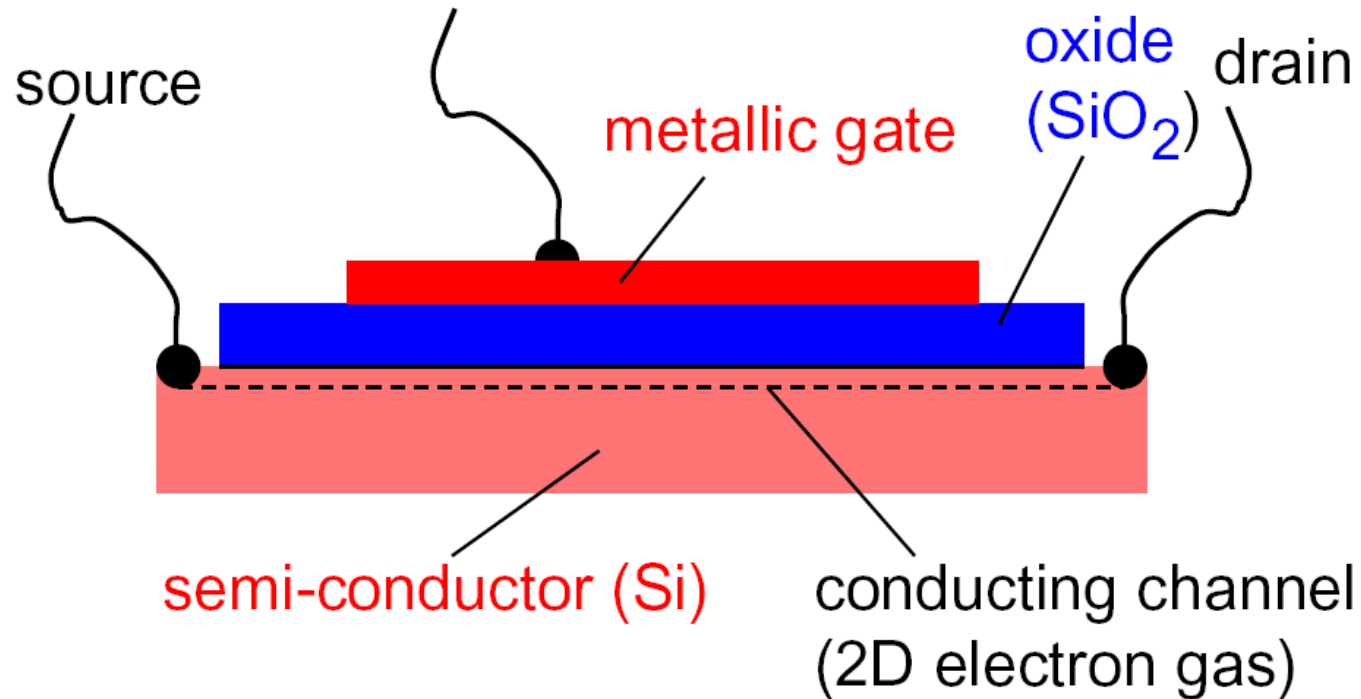
# What is an electron?

electron = object with **charge**  $-e$  and **spin**  $\frac{1}{2}$



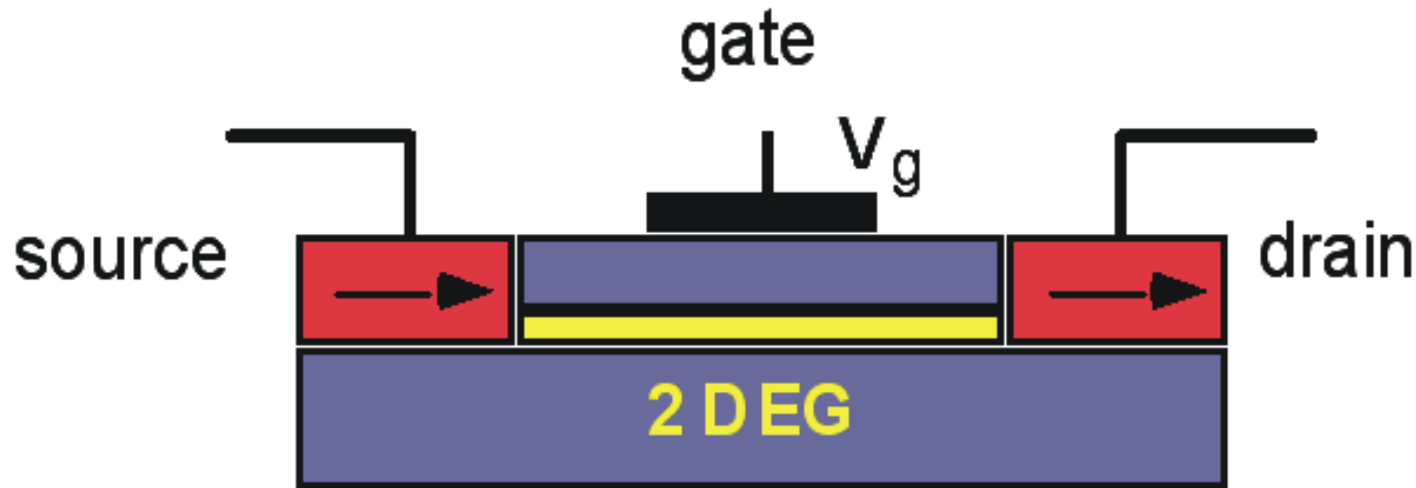
but so far, electronics has completely neglected the spin

# “Normal” transistor (MOSFET)



- gate voltage controls the current between source and drain
- used as modulator / amplifier or switch

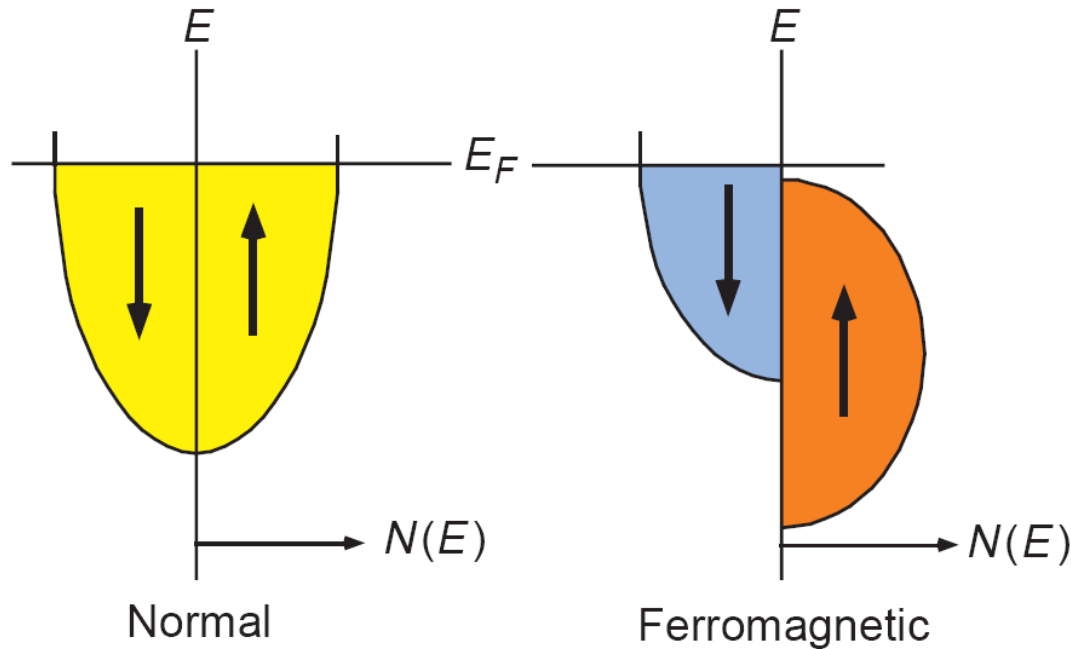
# Spin transistor



- source and drain are ferromagnetic materials
- gate voltage changes the spin of the electrons to control the current – **no external magnetic field**

# Spin-polarized transport

schematic DOS

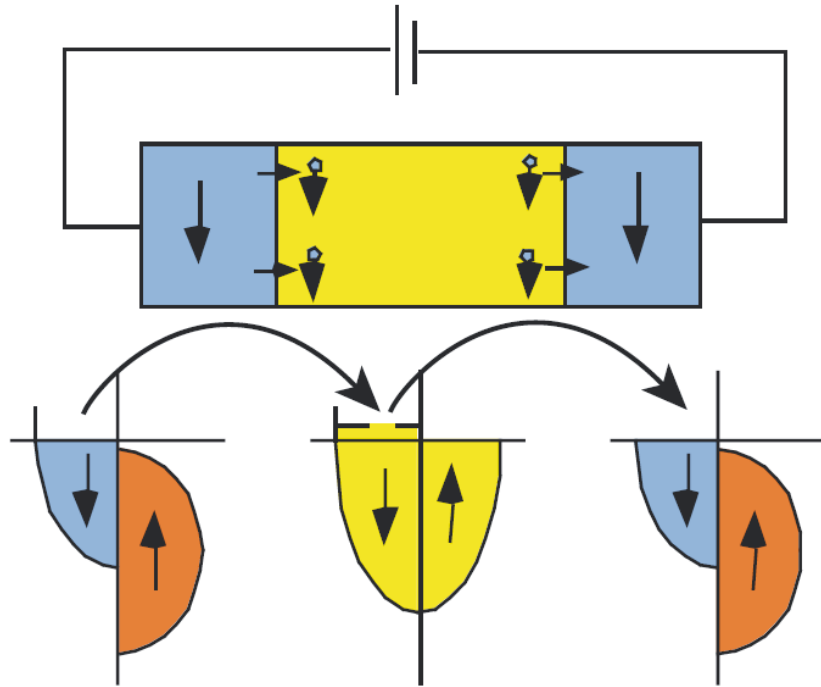


polarization:

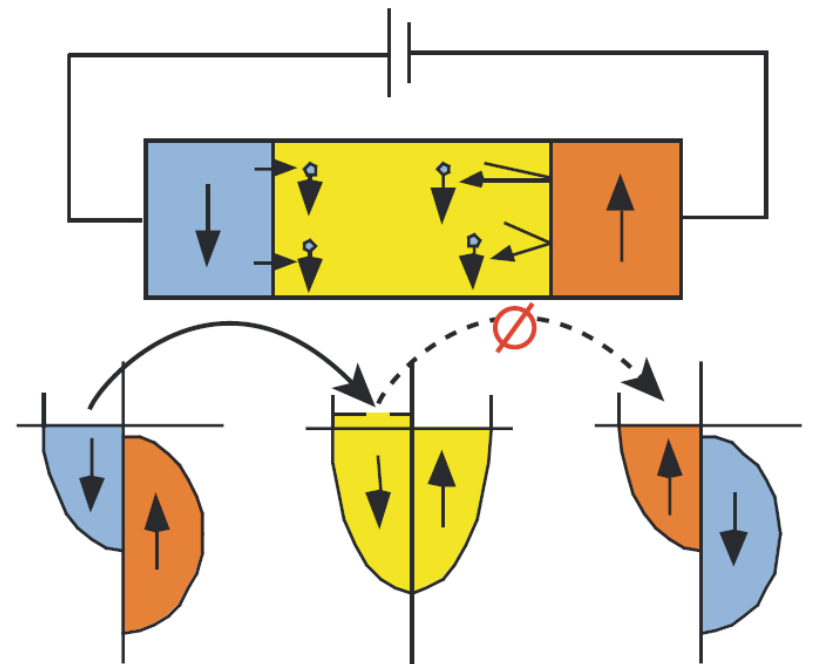
$$P = \frac{n_{\uparrow} - n_{\downarrow}}{n_{\uparrow} + n_{\downarrow}}$$

- imbalance of spin population at Fermi level leads naturally to spin-polarized transport
- commonly occurs in ferromagnetic metals (or alloys) with  $P$  up to 50 %

# Spin valve



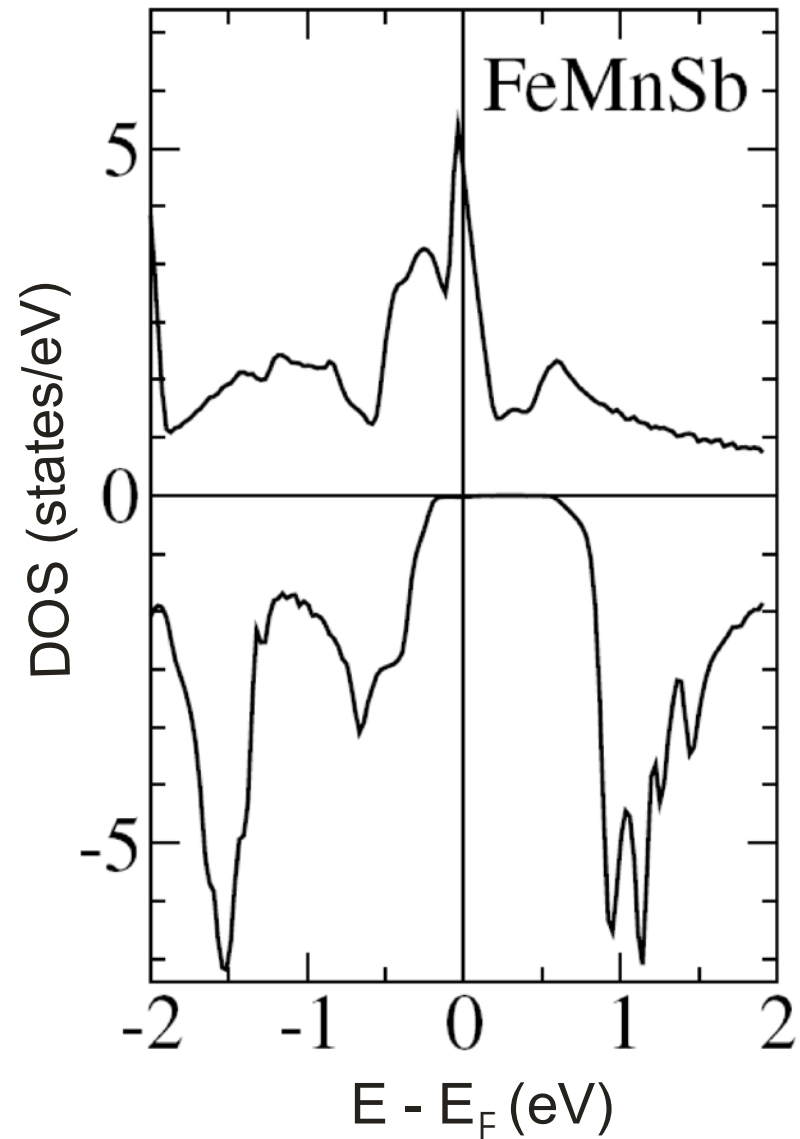
low resistance



high resistance

# New materials

- Heusler alloys have a (theoretical) spin polarization of 100 %
- e.g.  $\text{Co}_2\text{MnSi}$ ,  $\text{FeMnSb}$



Mavropoulos *et al.*  
J. Phys.: Condens. Matter **16** (2004)

# Spin-orbit coupling

- an electron moving in an electric field experiences a magnetic field in its rest frame of reference

$$\vec{B} \propto \vec{v} \times \vec{E}$$

- non-relativistic limit of Dirac equation leads to the Schrödinger equation with an additional term that describes spin-orbit coupling (SOC)

$$\hat{H}_{soc} = \frac{\hbar}{(2mc)^2} \nabla V \cdot (\hat{\sigma} \times \hat{k})$$

# Symmetry considerations

- from time reversal symmetry (Kramer's theorem)

$$E(\vec{k}, \uparrow) = E(-\vec{k}, \downarrow)$$

- from inversion symmetry

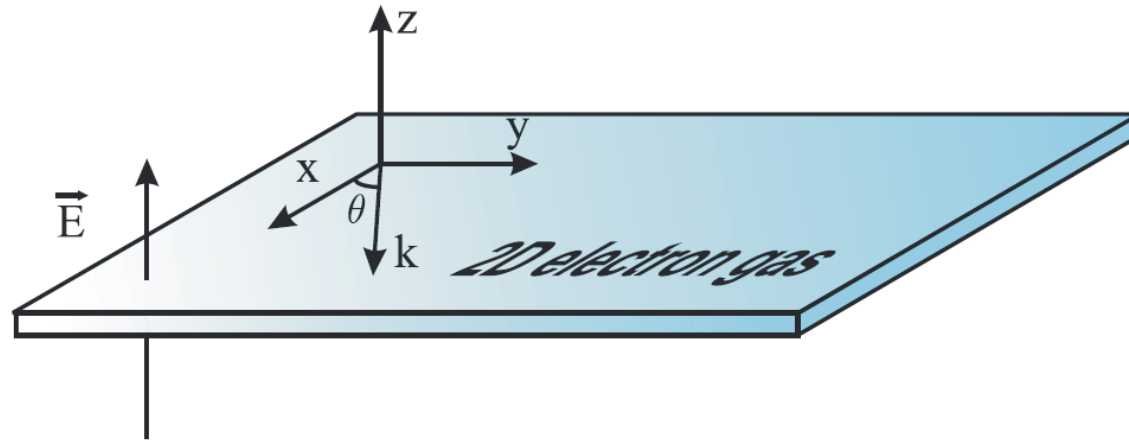
$$E(\vec{k}, \uparrow) = E(-\vec{k}, \uparrow)$$

- if both symmetries are fulfilled the bandstructure is doubly degenerated

$$E(\vec{k}, \uparrow) = E(\vec{k}, \downarrow)$$



# quasi-2D free electron gas



on the surface of a solid or in heterostructures the inversion symmetry is broken and SOC can lift the degeneracy of the states

# Rashba term

- spin-orbit Hamiltonian in 2D is so-called Rashba term

$$\begin{aligned} H_{SOC} &= \alpha (\vec{\sigma} \times \vec{k}) \vec{e}_z \\ &= \alpha (\sigma_x k_y - \sigma_y k_x) \end{aligned}$$

- coupling parameter  $\alpha$  can be controlled by the electric field

$$\alpha \propto \nabla V$$

# Ansatz

- full Hamiltonian of the 2DEG

$$H = \begin{pmatrix} \frac{1}{2m} (k_x^2 + k_y^2) & \alpha(k_y + ik_x) \\ \alpha(k_y - ik_x) & \frac{1}{2m} (k_x^2 + k_y^2) \end{pmatrix}$$

$$H|\psi\rangle = E|\psi\rangle$$

- commutator  $[k, H]_- = 0$

- wave function  $|\psi\rangle = e^{i(k_x x + k_y y)} (a|\uparrow\rangle + b|\downarrow\rangle)$

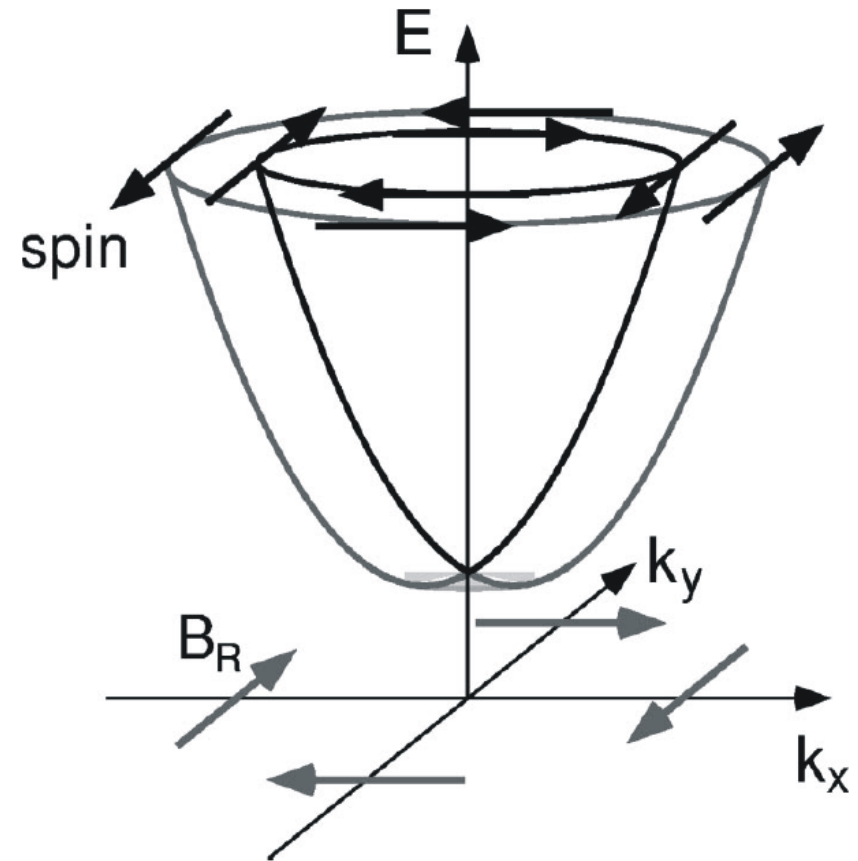
# Eigenvalues and eigenstates

$$E_1(k, \uparrow) = \frac{k^2}{2m} + \alpha|k|$$

$$E_2(k, \downarrow) = \frac{k^2}{2m} - \alpha|k|$$

$$\psi_1(\vec{r}) = \frac{1}{\sqrt{2}} e^{i\vec{k}\vec{r}} \left( |\uparrow\rangle + ie^{i\Theta} |\downarrow\rangle \right)$$

$$\psi_2(\vec{r}) = \frac{1}{\sqrt{2}} e^{i\vec{k}\vec{r}} \left( ie^{-i\Theta} |\uparrow\rangle + |\downarrow\rangle \right)$$



# Spin precession

- phase shift between two electrons with the same energy and different spin:

$$\Delta k = 2m\alpha$$

$$\Delta\theta = 2m\alpha L$$

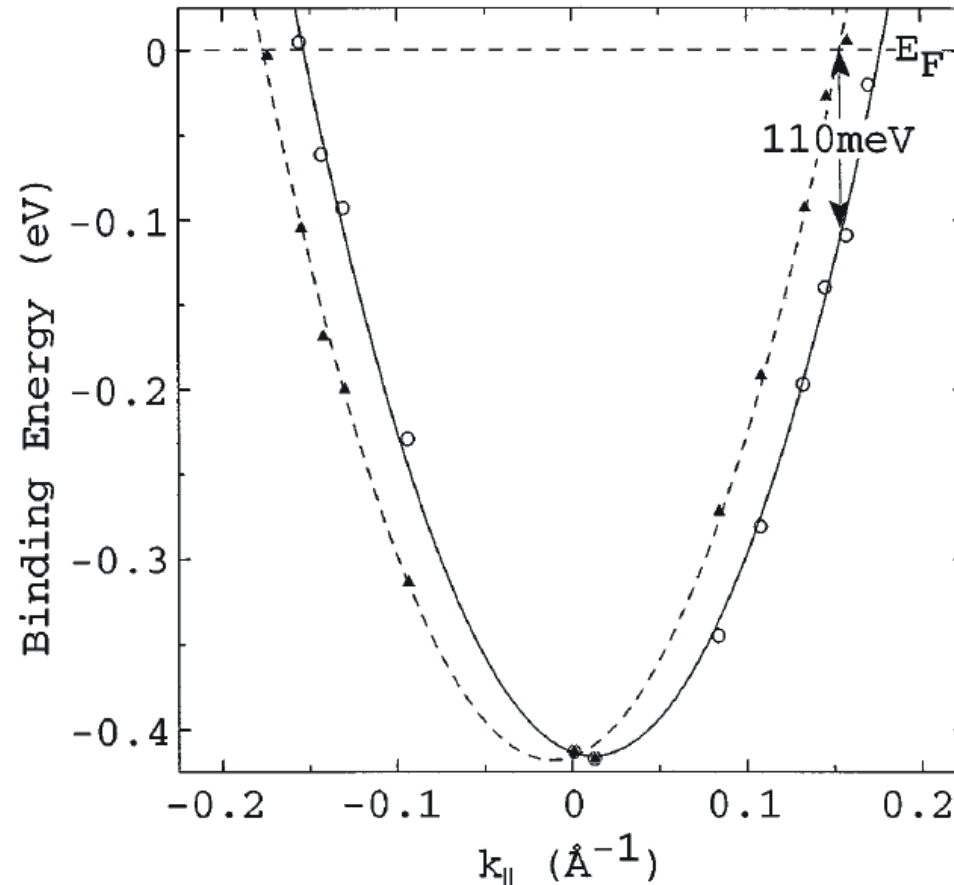
- spin rotates analogue to Larmor precession with frequency

$$\Omega(\vec{k}) = \alpha(\vec{k} \times \vec{n})$$

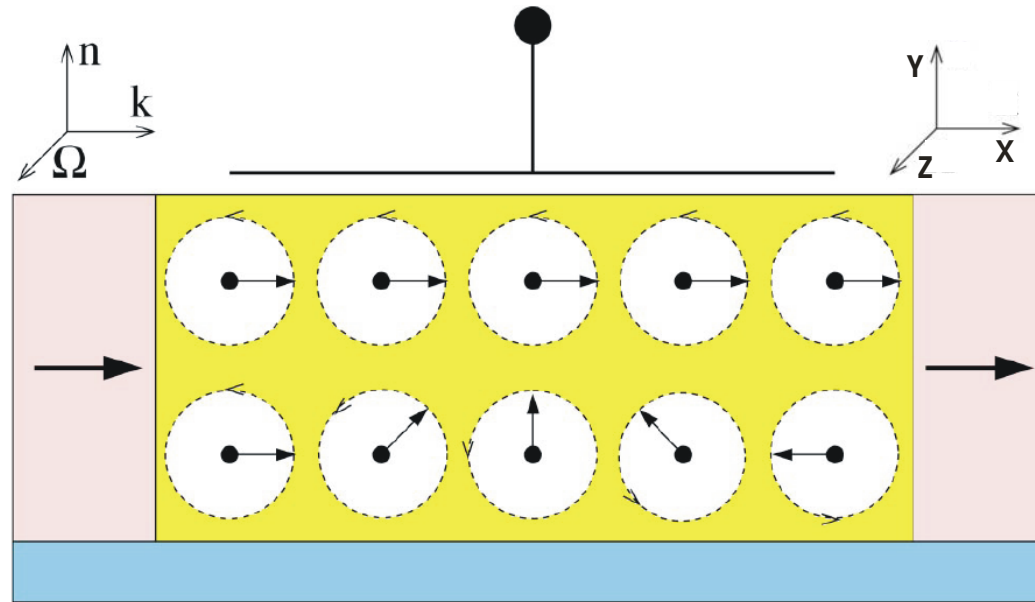
# Experimental results

spin splitting of Au(111) surface state

LaShell *et al.* PRL **77** (1996)



# Spin transistor revisited



x-polarized spin is a superposition of spin-up and spin-down eigenstates along the z-axis

$$\begin{pmatrix} 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \leftrightarrow \quad |x\rangle = |\uparrow\rangle + |\downarrow\rangle$$