

Topic 4: The Finite Potential Well

Outline:

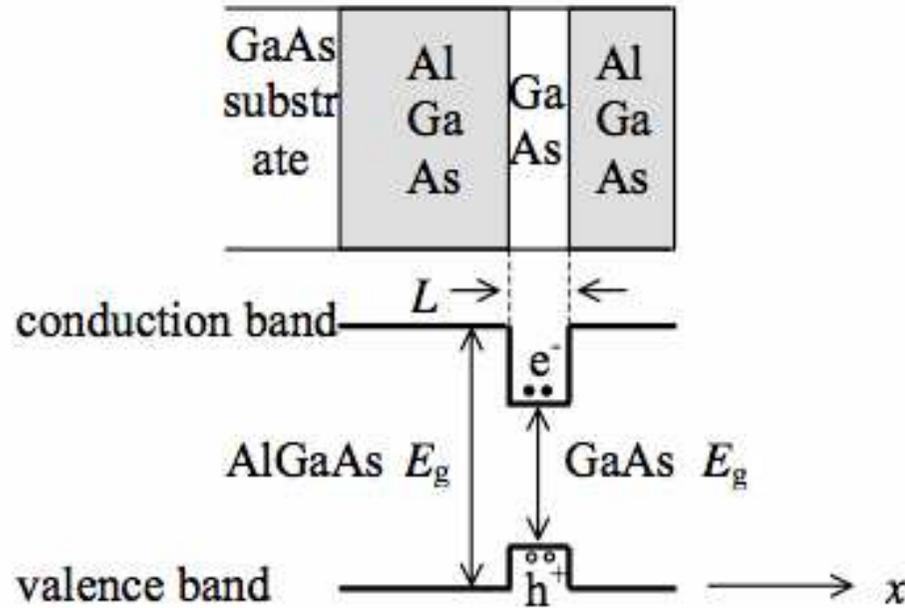
- The quantum well
- The finite potential well (FPW)
- Even parity solutions of the TISE in the FPW
- Odd parity solutions of the TISE in the FPW
- Tunnelling into classically forbidden regions
- Comparison with the IPW

The AlGaAs-GaAs Quantum Well

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Example of a potential well:

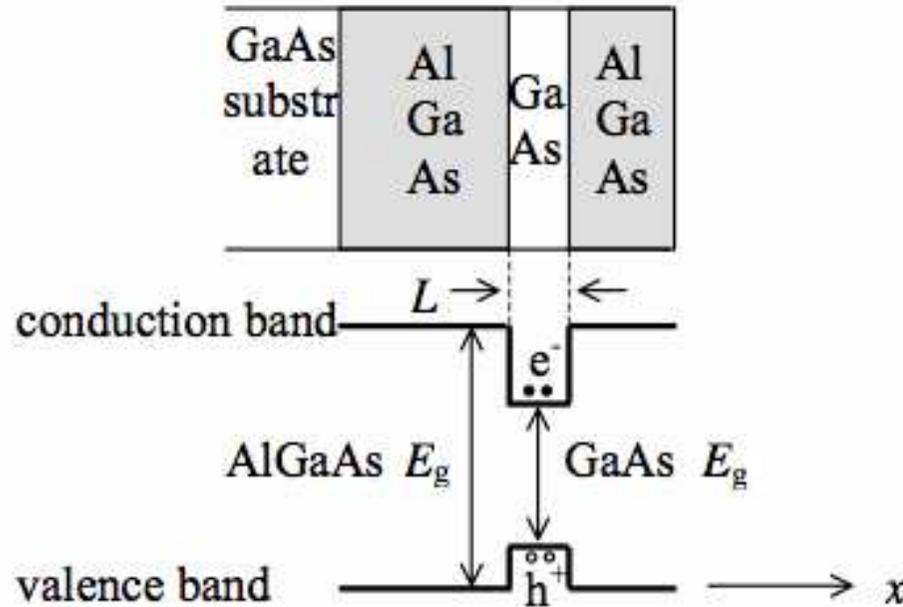
Sandwich of GaAs and AlGaAs layers



The AlGaAs-GaAs Quantum Well

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Constrained motion along the x -axis; free motion in the $y - z$ plane.

- $V_0 \rightarrow \infty$: 1-dimensional infinite potential well
- $V_0 < \infty$: 1-dimensional finite potential well

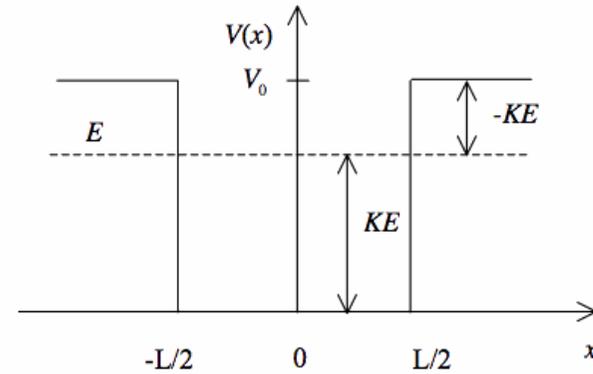
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Consider a particle in the potential

$$V(x) = \begin{cases} V_0 & x < -\frac{L}{2} \\ 0 & -\frac{L}{2} \leq x \leq \frac{L}{2} \\ V_0 & x > \frac{L}{2} \end{cases}$$

$$V_0 > 0$$

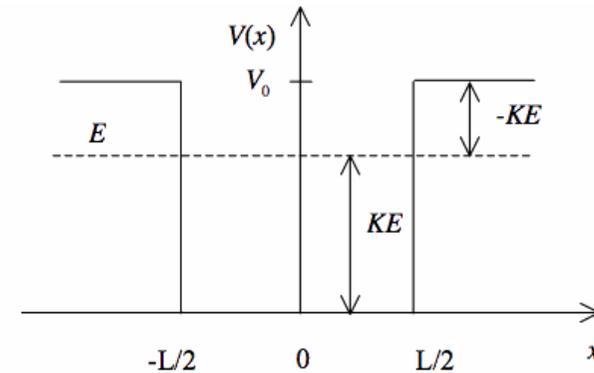


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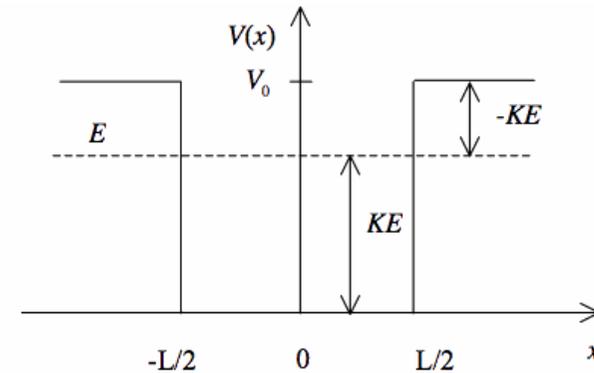
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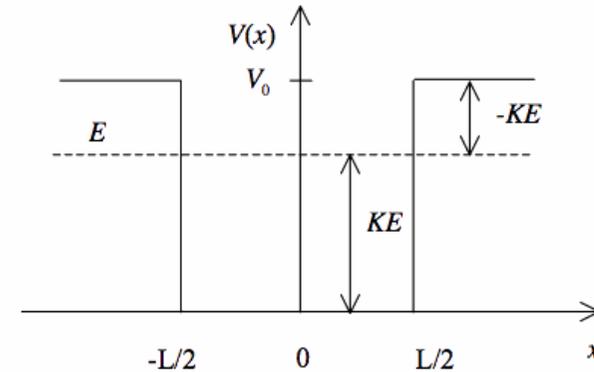
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Solve the TISE:

- $-\frac{L}{2} \leq x \leq \frac{L}{2}$ (Region I):

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) = E\psi(x) \Rightarrow \psi''(x) = -k^2\psi(x) \quad \boxed{k^2 = \frac{2mE}{\hbar^2} > 0}$$

solutions: $\psi_I(x) = A \sin kx + B \cos kx,$

(as for the IPW)

A, B – arbitrary constants

• $x > \frac{L}{2}$ (Region II):

Note: in region II $E = KE + PE = KE + V_0 < V_0 \Rightarrow KE < 0!$

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$$\Rightarrow \psi''(x) = \alpha^2 \psi(x)$$

$$\alpha^2 = \frac{2m(V_0 - E)}{\hbar^2} > 0$$

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C, D – arbitrary constants

\Rightarrow put $D = 0$, otherwise $\psi(x)$ not square integrable (blows up at large +ve x)

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• $x < -\frac{L}{2}$ (Region III):

$$\text{solutions } \psi_{III}(x) = Fe^{-\alpha x} + Ge^{\alpha x}$$

(like in region II)

F, G – arbitrary constants

\Rightarrow put $F = 0$, otherwise $\psi(x)$ not square integrable (blows up at large -ve x)

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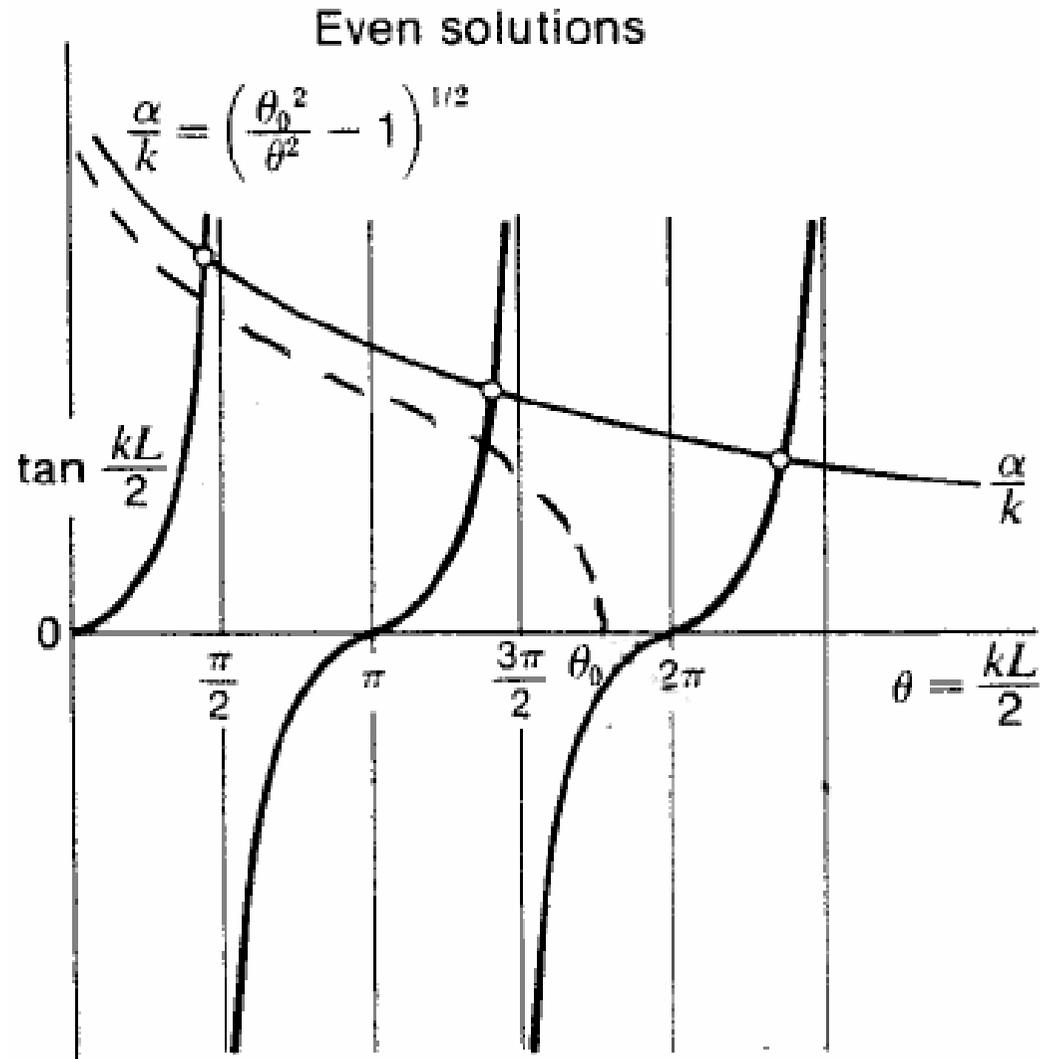
\Rightarrow RHS: $y(\theta) = \frac{\alpha}{k} = \sqrt{\frac{k_0^2}{k^2} - 1} = \sqrt{\frac{V_0 - E}{E}} = \sqrt{\frac{\theta_0^2}{\theta^2} - 1}$

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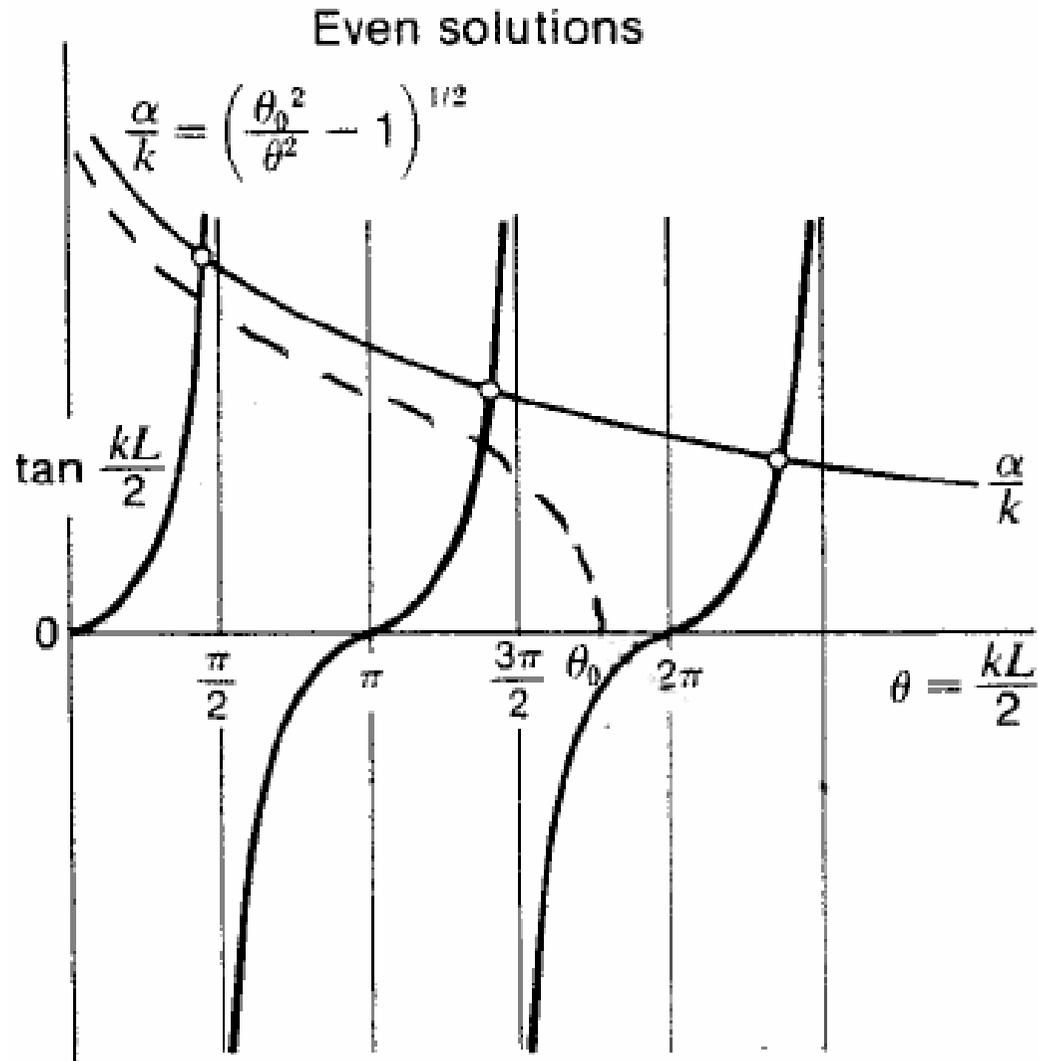
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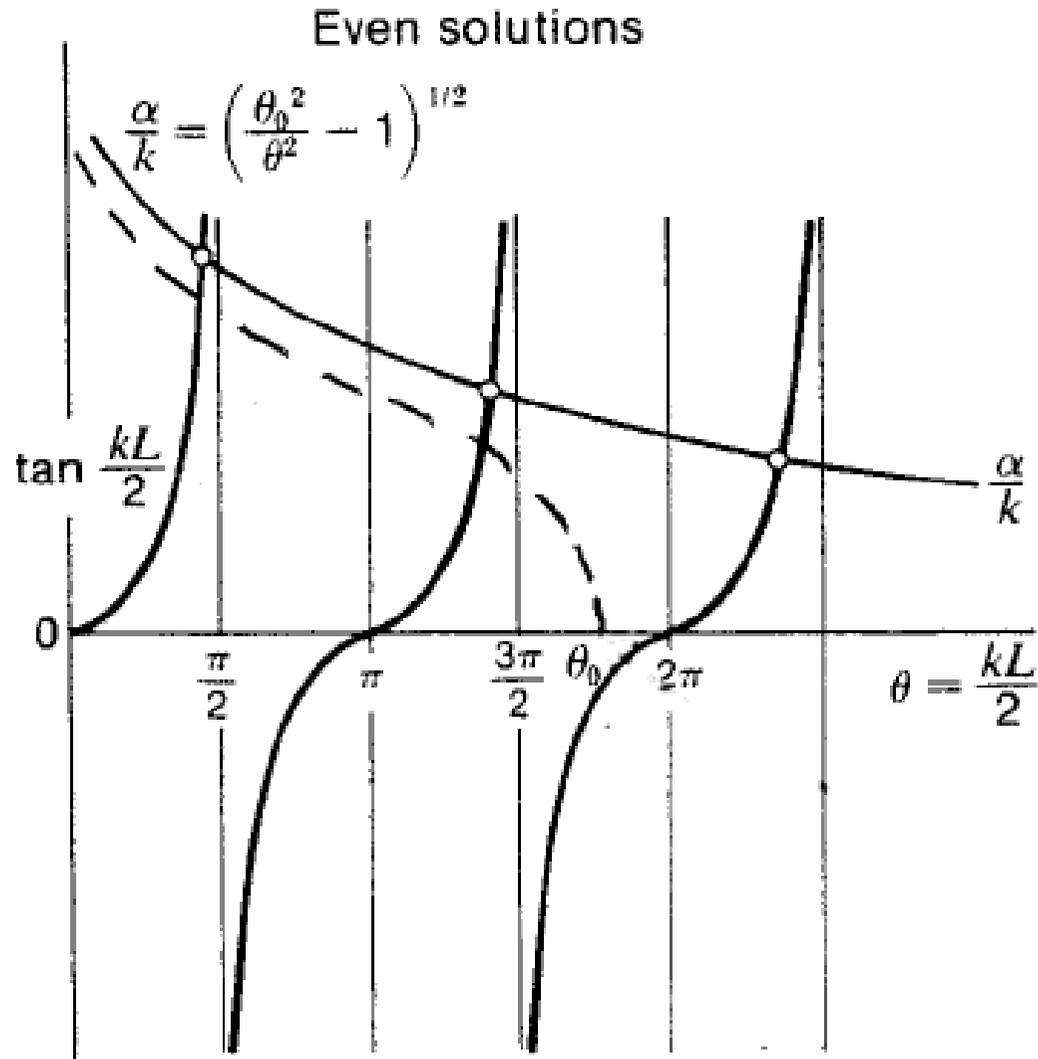
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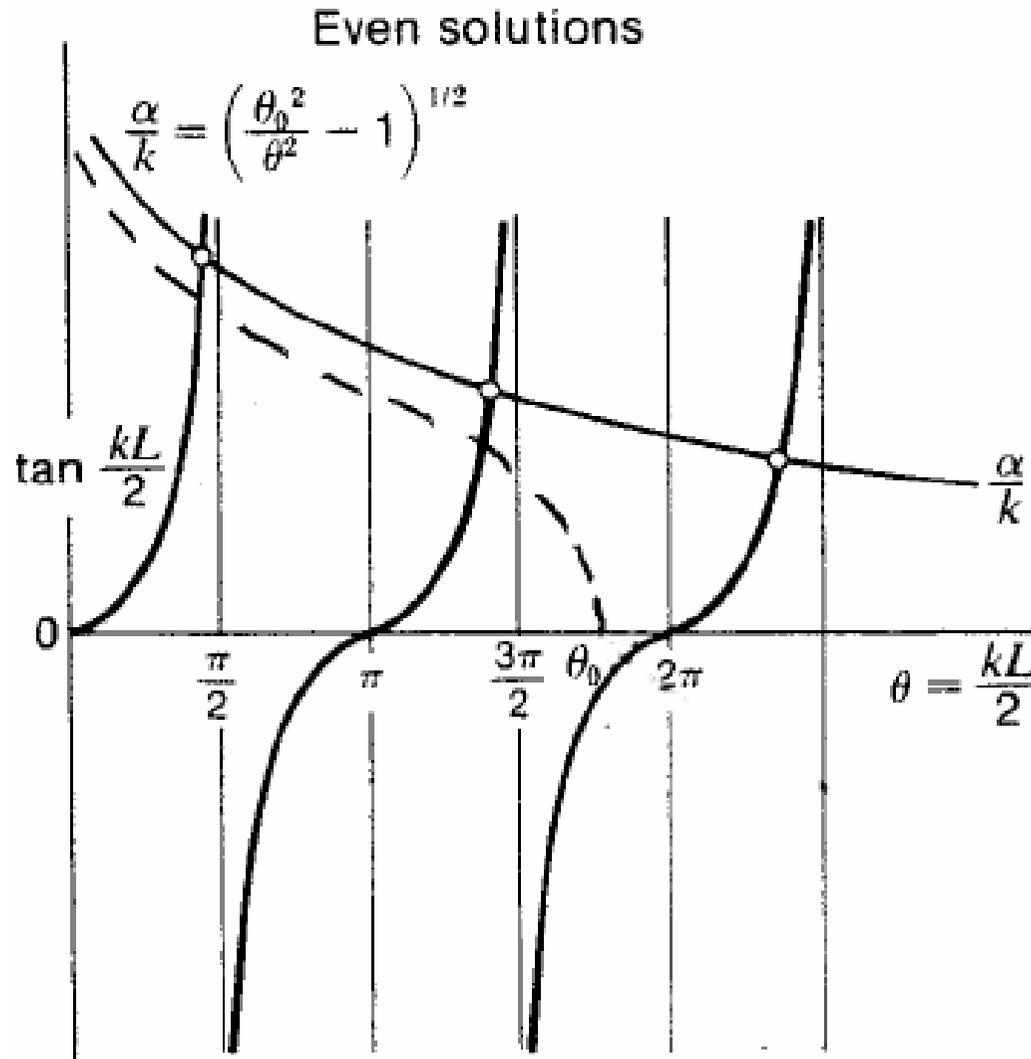
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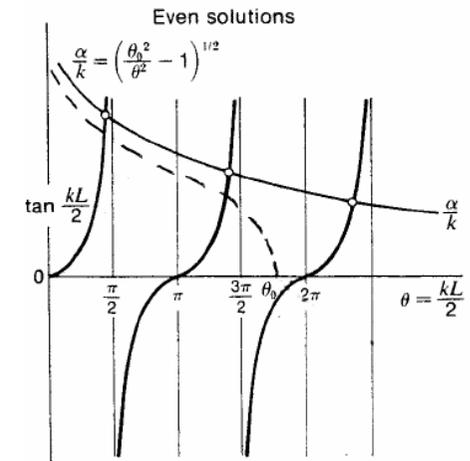


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\Rightarrow bound states with discrete (quantized) energy

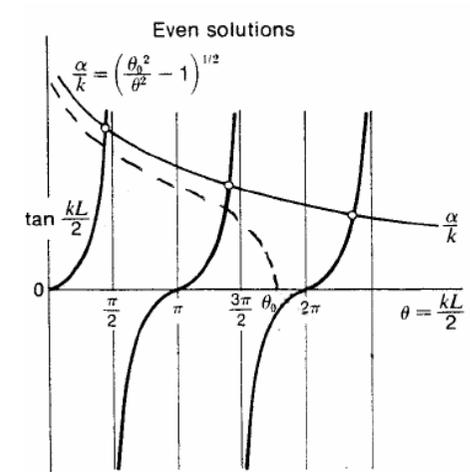
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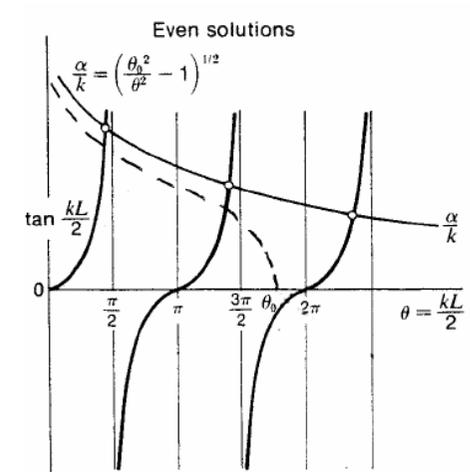
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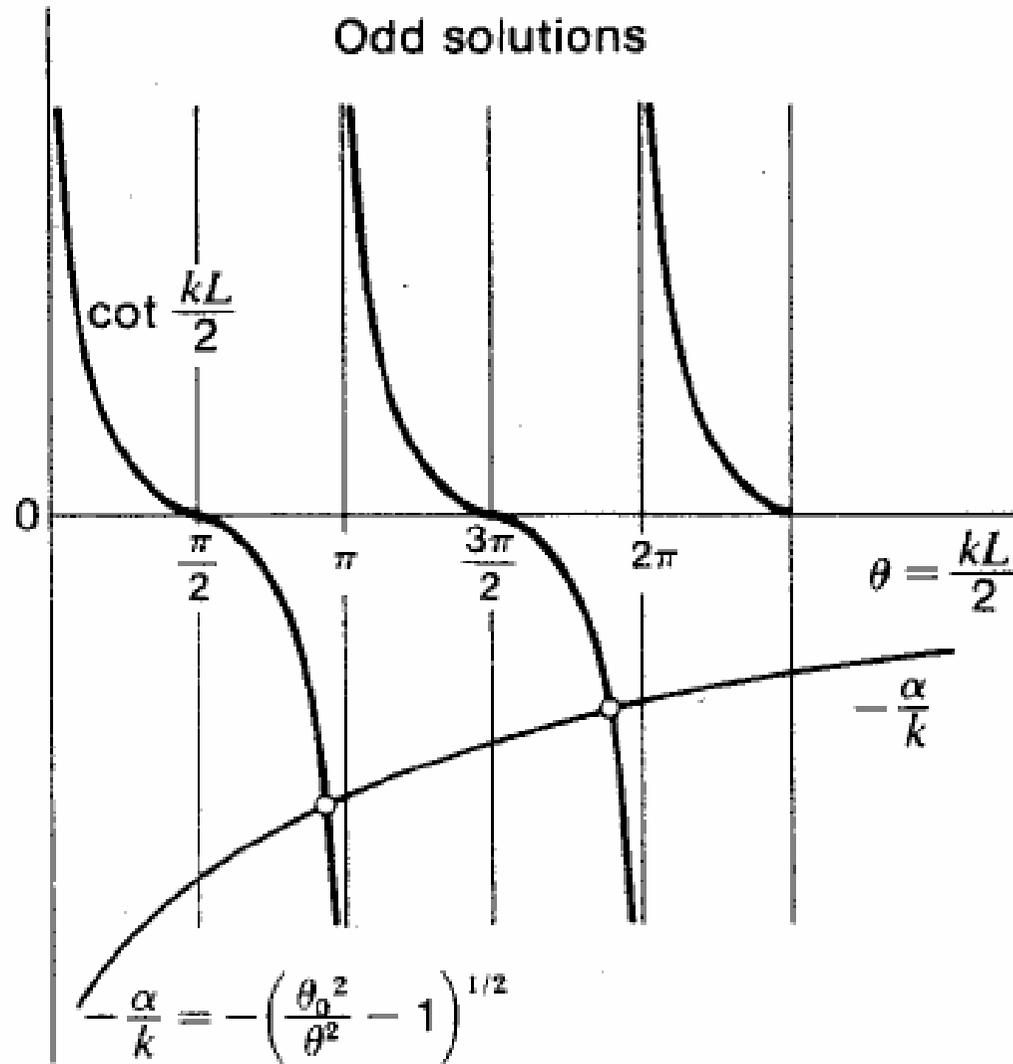
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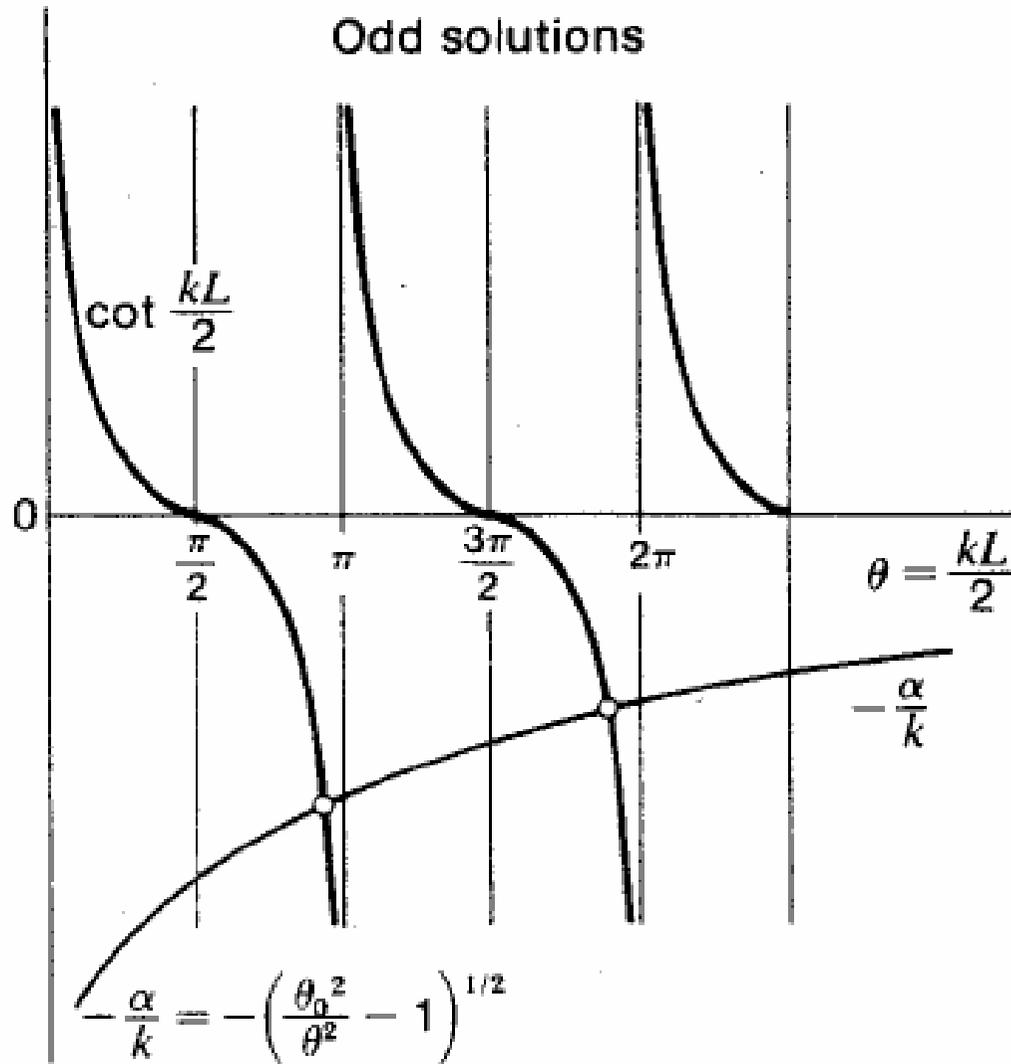
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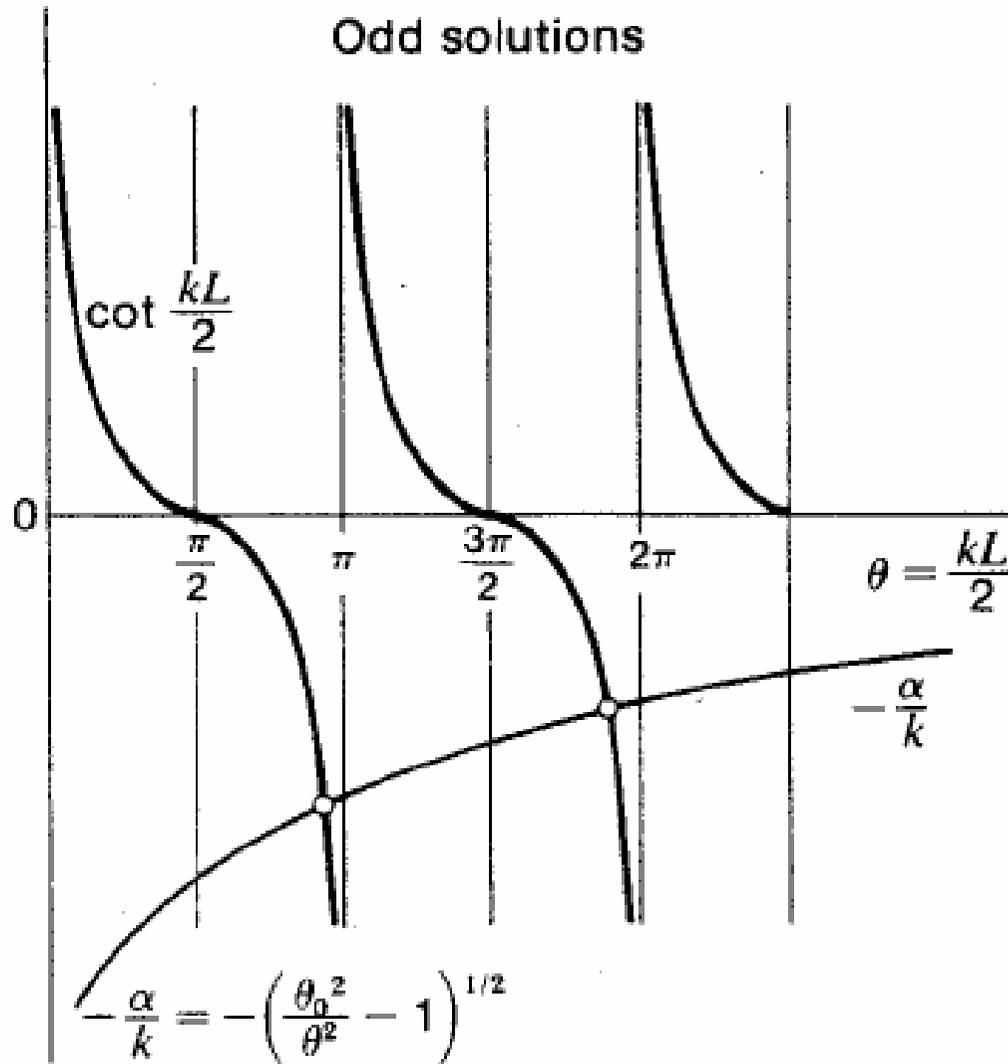
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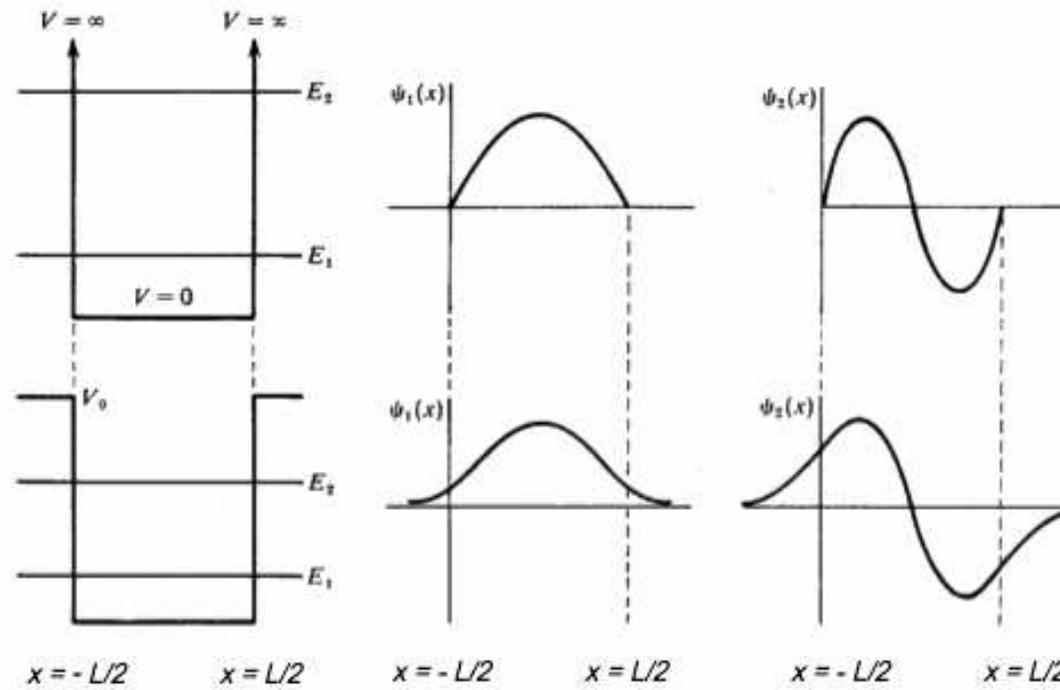
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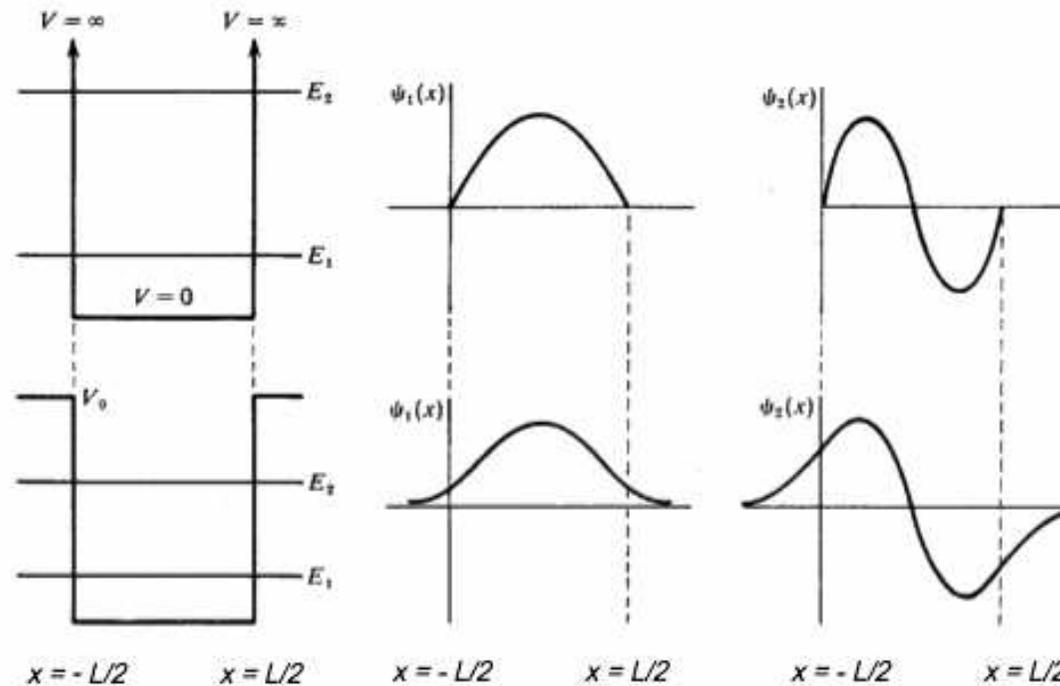


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Compare the FPW and the IPW



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Infinite well:

$\psi(x)$ confined to the well

$$k_n = \frac{n\pi}{L}$$

infinite tower of states

no unbound states

Finite well:

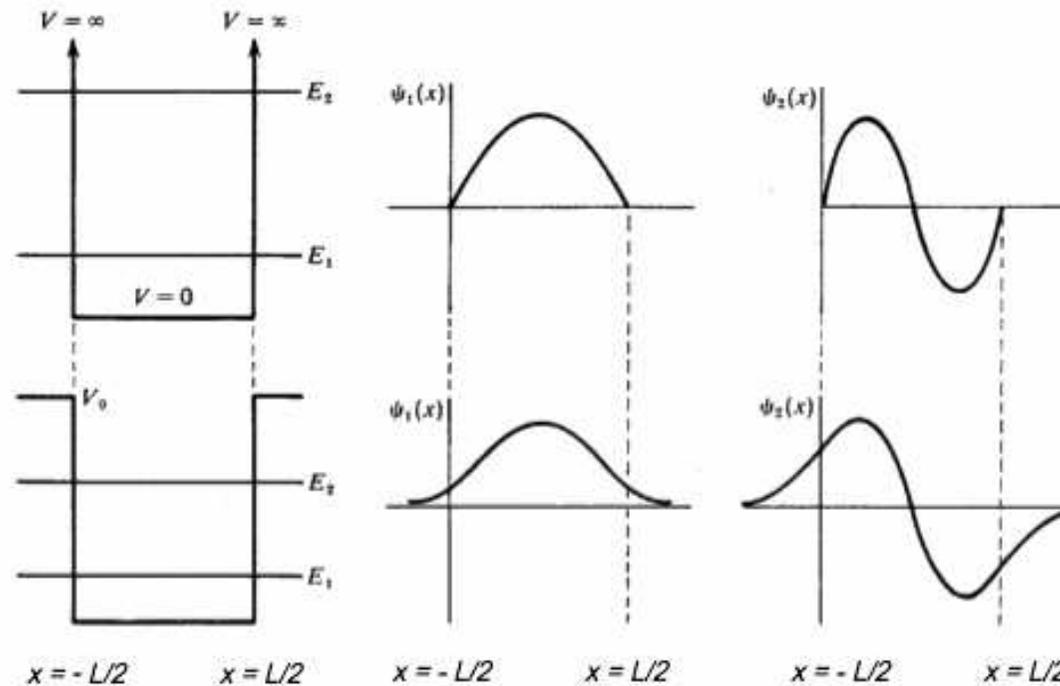
$\psi(x)$ spreads out beyond the well

k_n and energies lower

finite tower of states

unbound states when $E > V_0$

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The energy levels in the FPW are lower because the wavefunction spreads out (by penetrating the classically forbidden region) and therefore reduces its KE.

Quantum Tunnelling

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At $x > L/2$ the wavefn $\psi(x) \propto e^{-\alpha x}$;
at $x < -L/2$ the wavefn $\psi(x) \propto e^{\alpha x}$.

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- The depth of tunneling is determined by

α – the penetration depth

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 \Rightarrow Requiring a “reasonable behaviour” of the wavefunction leads to a (classically) “crazy” phenomenon of tunnelling

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- quantum states in symmetric potentials (w.r.t. reflections $x \rightarrow -x$) are either symmetric (i.e., even parity), with an even number of nodes, or else antisymmetric (i.e., odd parity), with an odd number of nodes

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- when $V = V(x)$, both bound and continuous states are stationary, i.e, the time-dependent wavefunctions are of the form

$$\Psi(x, t) = \psi(x) \exp\left(-\frac{i}{\hbar}Et\right)$$